

# Vlasov simulation of auroral processes

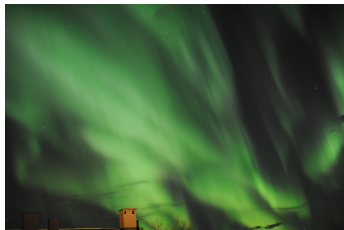
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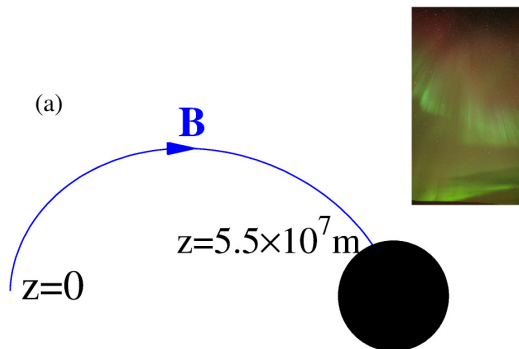
<sup>4</sup>Umeå University, Umeå, Sweden



## Simulation model

The distribution function is  $f_s(z, v_z, \mu)$  and the Vlasov equation

$$\frac{\partial f_s}{\partial t} + v_z \frac{\partial f_s}{\partial z} + \frac{1}{m_s} \left( q_s E - \mu \frac{\partial B}{\partial z} + m_s a_g \right) \frac{\partial f_s}{\partial v_z} = 0$$



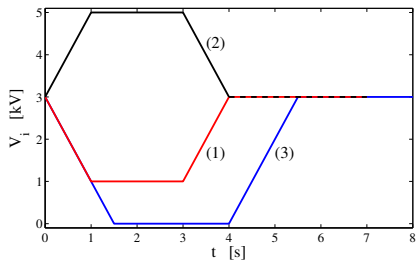
$$\mu = \frac{m_s v_{\perp}^2}{2B}$$

is a constant of motion.  
The electric field is given by

$$\frac{d}{dz} \left( \frac{B_S}{B} E \right) = \frac{\rho_l}{S\epsilon}$$

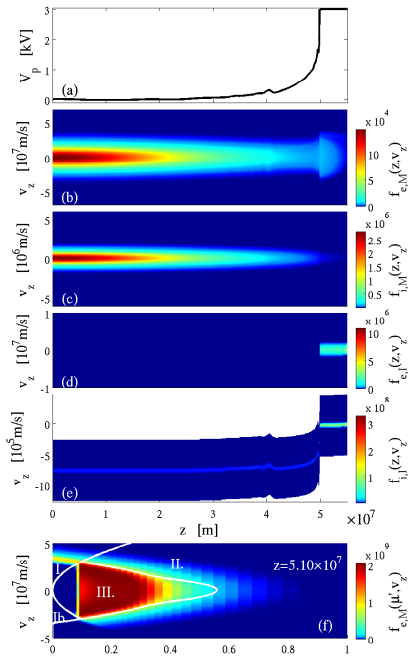
Computer code published  
by Gunell et al. (2013a)

# Changing the acceleration voltage during the simulation

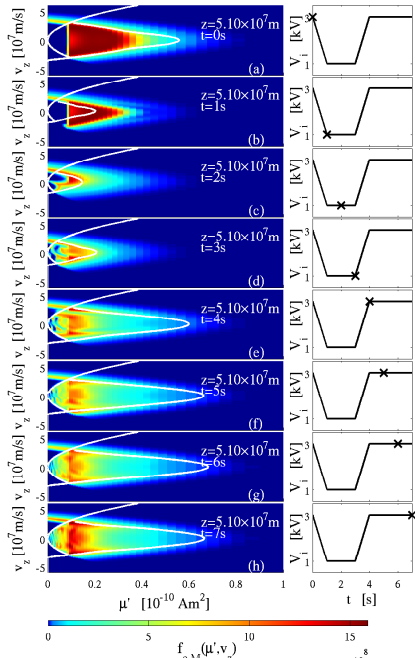


Above: Prescribed total voltage as a function of time.

Right: Initial state, common to all three experiments.

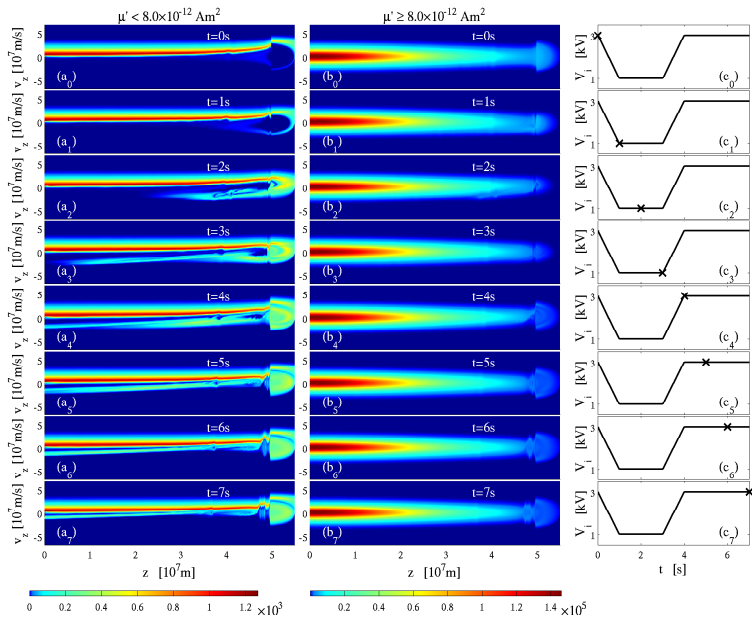


# Experiment 1: $f(\mu, v_z)$ at $z = 5.1 \times 10^7$ m

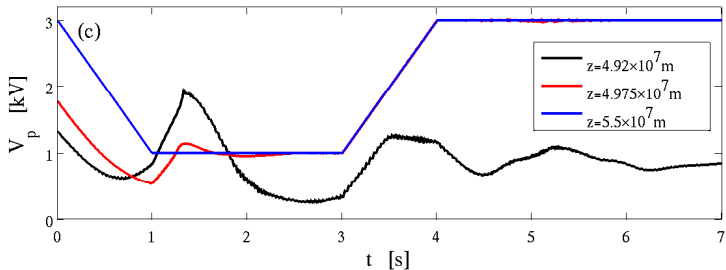
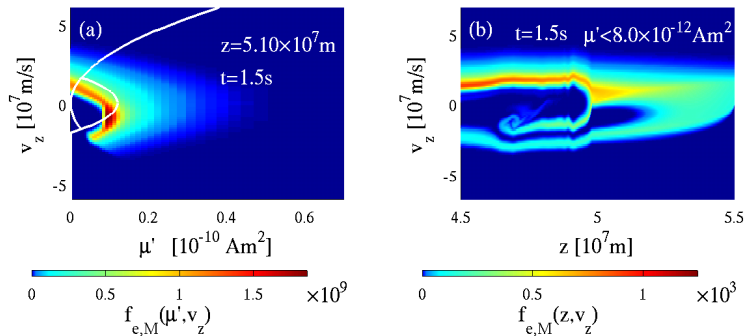


- ▶ Trapped electrons are released.
- ▶ New electrons get trapped.
- ▶ The distribution in the  $\mu$  direction changes.

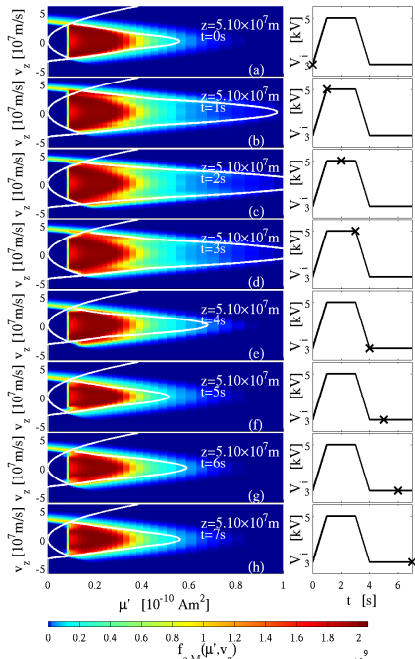
# Experiment 1: $f(z, v_z)$



# Experiment 1: fluctuations

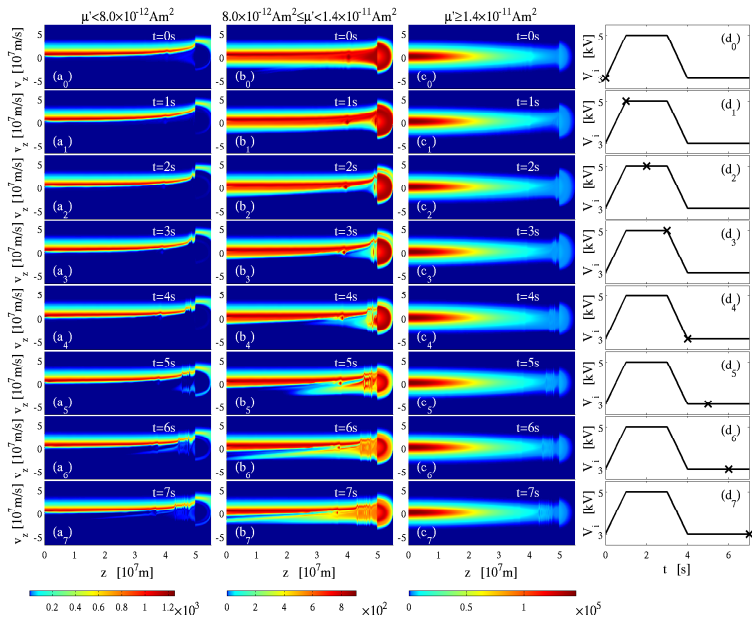


## Experiment 2: $f(\mu, v_z)$ at $z = 5.1 \times 10^7$ m



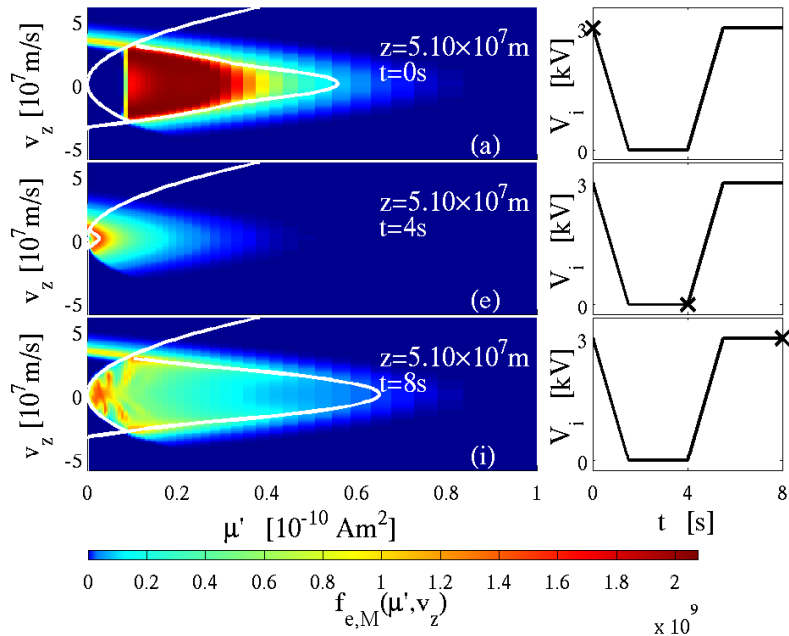
- ▶ New electrons are trapped at large  $\mu$  values.
- ▶ Then these are lost again.
- ▶ The final state resembles the initial, although there is new structure at intermediate  $\mu$  values.

# Experiment 2: $f(z, v_z)$

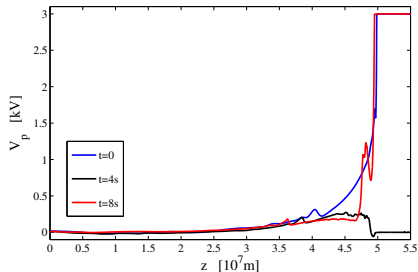
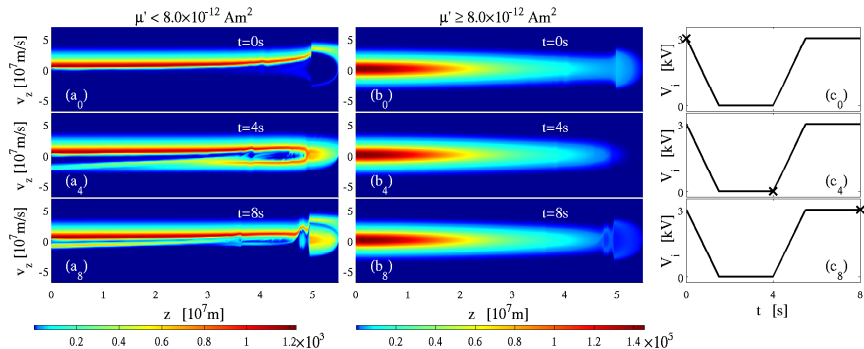




# Experiment 3: $f(\mu, v_z)$ at $z = 5.1 \times 10^7$ m

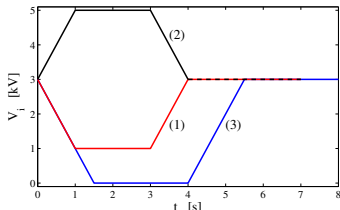
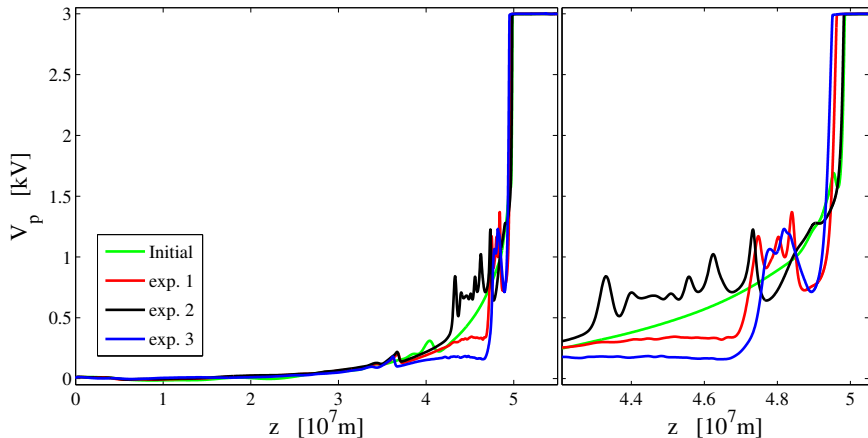


# Experiment 3: $f(z, v_z)$



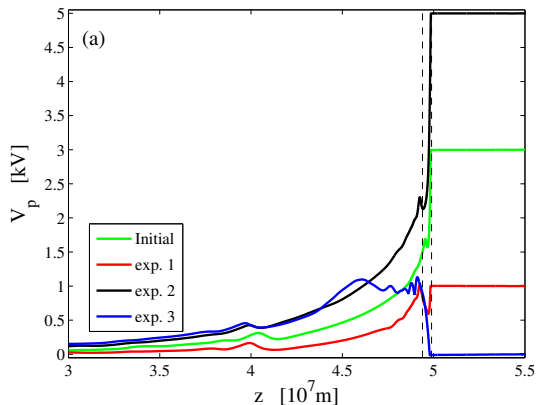
- ▶ The double layer changes polarity when the voltage is low.
- ▶ At  $t = 8$  s, a new potential peak traps some electrons.

## Final state

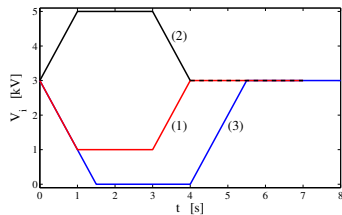
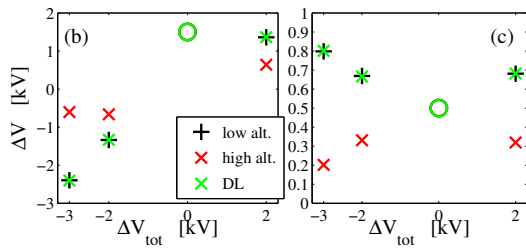


- ▶ The potential structure changes on the low potential side of the double layer.
- ▶ The double layer moved for  $\Delta V < 0$ .

# Where does $\Delta V$ go?

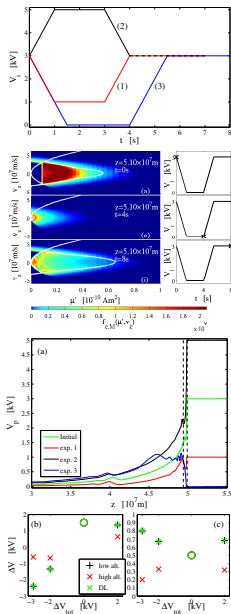


Most of  $\Delta V$  is assumed by the double layer, c.f. observations by Forsyth et al. (2012).



# Summary

- ▶ Trapping and release processes can be followed in simulations.
- ▶ The history of the acceleration voltage affects the distribution function of the trapped electrons.
- ▶ Most of the change in the acceleration voltage is assumed by the double layer.
- ▶ For a quickly decreasing voltage the double layer changes polarity.
- ▶ Phase space holes are created in these numerical experiments.
- ▶ The double layer position exhibits hysteresis phenomena.



## Bibliography

Forsyth, C., Fazakerley, A. N., Walsh, A. P., Watt, C. E. J., Garza, K. J., Owen, C. J., Constantinescu, D., Dandouras, I., Fornaçon, K.-H., Lucek, E., Marklund, G. T., Sadeghi, S. S., Khotyaintsev, Y., Masson, A., Doss, N., Dec. 2012. Temporal evolution and electric potential structure of the auroral acceleration region from multispacecraft measurements. *Journal of Geophysical Research (Space Physics)* 117 (A16), 12203.

URL <http://dx.doi.org/10.1029/2012JA017655>

Gunell, H., De Keyser, J., Gamby, E., Mann, I., 2013a. Vlasov simulations of parallel potential drops. *Annales Geophysicae* 31, 1227–1240.

URL <http://dx.doi.org/10.5194/angeo-31-1227-2013>

Gunell, H., De Keyser, J., Mann, I., 2013b. Numerical and laboratory simulations of auroral acceleration. *Physics of Plasmas* 20, 102901.

URL <http://www.herbertgunell.se/pdfpapers/GDM13.pdf>

# Vlasov simulations – technical details – $\epsilon_r$

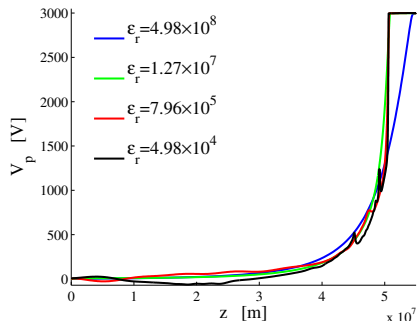
## Computational advantages

$\omega_p \sim 1/\sqrt{\epsilon_r} \Rightarrow$  lets us use a longer time step

$\lambda_D \sim \sqrt{\epsilon_r} \Rightarrow$  lets us use a coarser grid

## Disadvantage

Suppression of small scale and high frequency phenomena



	Magneto- sphere	Iono- sphere
$z$	0	$5.5 \cdot 10^7 \text{ m}$
$B$	$0.086 \mu\text{T}$	$56 \mu\text{T}$
$k_B T_e$	500 eV	1 eV
$k_B T_{H^+}$	2500 eV	1 eV
$n$	$3 \cdot 10^5 \text{ m}^{-3}$	$1 \cdot 10^9 \text{ m}^{-3}$